



Large Fluctuations and Fixation in Evolutionary Games with Non-Vanishing Selection

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- **Intro to Evolutionary game theory**
 - Basic notions & replicator Dynamics
 - Stochastic Dynamics, Evolutionary stability & Fixation
- **Large fluctuations & WKB theory in evolutionary games**
 - WKB theory in anti-coordination games
 - General WKB treatment & Results
 - WKB calculation of the fixation probability in coordination games
 - Comparison with diffusion approximations (Fokker-Planck)
- **Outlook & Conclusion**

Need for an accurate theoretical approach to describe large-fluctuation-induced phenomena (stochastic fluctuations and nonvanishing selection), diffusion approx don't work here

What is Evolutionary Game Theory about?

What is Evolutionary Game Theory about?

- Description of complex phenomena in behavioural science and population dynamics (e.g. in ecology, genetics, economics, ...)
- Dynamical version of *classic (rational) game theory*

Some of the founders & pioneers:

- John von Neumann & Oskar Morgenstern (1944), “Theory of games and economic behavior”
- John Nash (1994 Nobel prize in Economics) → **Nash equilibrium**
- John Maynard Smith, “Evolution and the Theory of Games” (1972) → **Evolutionary stability**

Some reference books:

- J. Hofbauer & K. Sigmund, “Evolutionary Games and Population Dynamics” (1998)
- M. Nowak, “Evolutionary Dynamics” (2006)
- J. Maynard Smith, “Evolution and the Theory of Games” (1972)

Some Basics of Game Theory

Initially, “game theory” was a branch of social sciences and applied maths (von Neumann & Morgenstern, 1944). Goal: find optimal strategies (“utility function”).

Evolutionary Game Theory (EGT): different approach where utility function (game’s payoff) is the **reproductive fitness** \Rightarrow successful strategies spread at the expenses of the others (Maynard Smith & Price, 1973).

New aspects and interpretations:

- 1 Strategies and their frequencies become *population species* and their *densities*
- 2 Dynamics is naturally implemented in EGT

The Replicator Dynamics

Traditional EGT setting: large and unstructured populations with pairwise interactions.

At *mean-field* level, the dynamics is described by the replicator equations for the density x_i of type $i = 1, \dots, S$ in the population:

$$\dot{x}_i = x_i(\Pi_i - \bar{\Pi}),$$

where Π_i : average payoff (here = fitness) of an individual of species i
 $\bar{\Pi}$: mean payoff averaged over the entire population

Common choice, with a payoff matrix \mathcal{P} : $\Pi_i = (\mathcal{P}\mathbf{x})_i$ linear function of $\mathbf{x} = (x_1, \dots, x_i, \dots, x_S)$, $\bar{\Pi} = \mathbf{x} \cdot \mathcal{P}\mathbf{x}$

Important case: 2×2 games with 2 species/strategies (A and B)

vs	A	B
A	a	b
B	c	d

A vs A gets a and B vs B gets d ; A vs B gets b , while B gets c

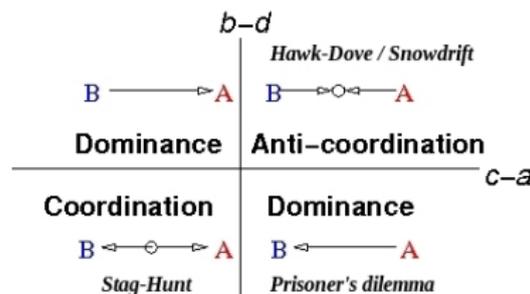
Replicator Dynamics for 2×2 Games

Population comprised of a density x and $1 - x$ of A and B , resp.
 Thus, $\Pi_A = ax + b(1 - x)$, $\Pi_B = cx + d(1 - x)$ and $\dot{\Pi} = x\Pi_A + (1 - x)\Pi_B$

$$\dot{x} = x(1-x)[(a-b-c+d)x + b-d] \Rightarrow$$

$$x^* = \frac{d-b}{a-b-c+d} \quad (\text{Interior fixed point})$$

- **Dominance** $(a - c)(d - b) \leq 0$: A dominates over B when $a \geq c$ & $b \geq d$. B dominates over A when $c \geq a$ & $d \geq b$
- **Coordination (bistability)**: When $a > c$ and $d > b$, the absorbing states $x = 0$ and $x = 1$ are stable and separated by x^* (unstable)
- **Anti-coordination (coexistence)**: When $c > a$ and $b > d$, x^* is stable while $x = 0$ and $x = 1$ are unstable, hence A and B coexist



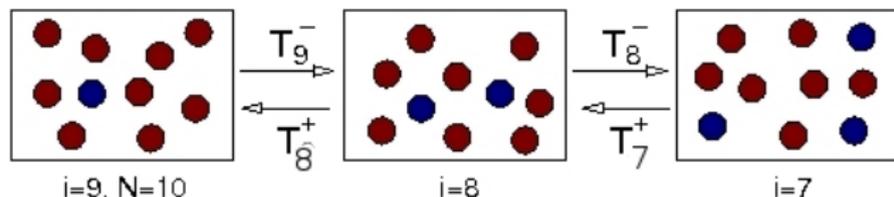
- **Neutrality**: When $a = c$ and $b = d$, there is neutral stability for all values of x

Stochastic Dynamics & Moran Process

Evolutionary dynamics involves a *finite number of discrete individuals*
⇒ stochastic rules given by the frequency-dependent **Moran process**

2×2 games: Markov birth-death process of i individuals of species A and $N - i$ of species B (total size N is conserved).

- At each time step, randomly pick 2 individuals
- 1 individual selected for reproduction and the other for death. The offspring replaces the deceased. N remains constant
- “Interaction” according to the payoff matrix, i.e. reproduction and death rates depend on the individuals’ fitnesses f_A and f_B .
- Transition $i \rightarrow i + 1$ (birth of a A and death of a B) with rate T_i^+ , while the transition $i \rightarrow i - 1$ (birth of B and death of a A) occurs with rate T_i^- . T_i^\pm **are functions of f_A and f_B**



The probability $P_i(t)$ of having i individuals of species **A** at time t obeys the master equation:

$$\frac{d}{dt}P_i(t) = T_{i-1}^+ P_{i-1}(t) + T_{i+1}^- P_{i+1}(t) - [T_i^+ + T_i^-]P_i$$

$i = 0$ (i.e. all **B**'s) and $i = N$ (i.e. all **A**'s) are **absorbing states** \Rightarrow
 $i \in [0, N]$ and $T_0^\pm = T_N^\pm = 0$

For the frequency-dependent Moran Process (fMP):

- Fitnesses of **A** and **B** given by f_A and f_B , resp.
 2 contributions: baseline (neutral) contribution + *selection* \Rightarrow
 $f_A = 1 - w + w\Pi_A$ and $f_B = 1 - w + w\Pi_B$.
 Strength of selection measure by $0 \leq w \leq 1$:
 $w = 0 \rightarrow$ neutrality, $w = 1 \rightarrow$ only selection
- $T_i^\pm = \chi_i^\pm(f_A, f_B)$
- $\chi_i^+ = \frac{f_A}{(i/N)f_A + (1-i/N)f_B}$ and $\chi_i^- = \frac{f_B}{(i/N)f_A + (1-i/N)f_B} \rightarrow \frac{f_B}{xf_A + (1-x)f_B}$

Markov chain with absorbing boundaries \Rightarrow unavoidable **fixation**,
 with system ending with all **A**'s ($i = N$) or all **B**'s ($i = 0$)

Stochastic fluctuations alter the predictions of the replicator equations

Evolutionary Stability & Fixation

- 1 Fixation: possibility for a few mutants to take over the entire population
- 2 There is *evolutionary stability* when the population *B*'s is proof against invasion from mutants *A*'s

Starting with i mutants of type *A*, what is the probability ϕ_i^A of ending with all *A*'s ($i = N$)? How long does it take? Dependence on w ?

In the neutral case ($w = 0$), $\phi_i^A = i/N \Rightarrow$

State with all *B*'s evolutionary stable if selection opposes replacement by *A* mutants *A*, i.e. if $\phi_i^A < i/N$

2×2 evolutionary games are formulated as 1D single-step birth-death processes and thus (formally) solvable:

- $\phi_i^A = \frac{1 + \sum_{k=1}^{i-1} \prod_{l=1}^k \gamma_l}{1 + \sum_{k=1}^{N-1} \prod_{l=1}^k \gamma_l}$, with $\gamma_l = T_l^- / T_l^+ = \chi_l^- / \chi_l^+$
- Unconditional fixation time:
$$\tau_i = -\tau_1 \sum_{k=i}^{N-1} \prod_{m=1}^k \gamma_m + \sum_{k=i}^{N-1} \sum_{l=1}^k \frac{1}{T_l^+} \prod_{m=l+1}^k \gamma_m,$$

with $\tau_1 = \phi_1^A \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{1}{T_l^+} \prod_{m=l+1}^k \gamma_m$

Large Fluctuations & WKB-based Theory

- 1 Exact expressions: difficult to generalise and analyse
- 2 Common approach: Fokker-Planck approximation (FPA) → good only for weak selection (diffusive dynamics: tractable)
- 3 Evolutionary dynamics: generally combination of random fluctuations and non-vanishing selection → Other approach is needed

When there is *metastability* fixation is reached following an “optimal path” obtained by a WKB theory

- ACG: WKB analysis ⇒ quasi-stationary distribution (QSD), probability and mean times of fixation (MFTs)
- CG: WKB calculation of the fixation probability

Anti-coordination Games & WKB Theory (I)

In ACGs ($c > a, b > d$), after relaxation time t_r , the system converges to the **metastable state** $n_* = Nx^*$. The latter has a very long mean time of decay, τ , that coincides with the (unconditional) MFT

WKB treatment requires: (1) $\tau \gg t_r$, (2) N and $Nx^* \gg 1$, (3) transition rates of order $\mathcal{O}(1)$ away from the absorbing boundaries

Idea: At time $t \gg t_r$, $P_i(t) \simeq \pi_i e^{-t/\tau}$ for $1 \leq i \leq N-1$ and $P_0(t) \simeq \phi(1 - e^{-t/\tau})$, $P_N(t) \simeq (1 - \phi)(1 - e^{-t/\tau})$

From fluxes of probability into the absorbing states:

- Unconditional MFT: $\tau = [T_1^- \pi_1 + T_{N-1}^+ \pi_{N-1}]^{-1}$
- Conditional MFTs: $\tau^A = [T_{N-1}^+ \pi_{N-1}]^{-1}$ and $\tau^B = [T_1^+ \pi_1]^{-1}$
- Fixation probability: $\phi^B = 1 - \phi^A = \phi = T_1^- \pi_1 \tau$

This requires the full QSD π_i . Assuming π_i/τ negligible, the *quasi-stationary master equation* (QSME)

$$0 = T_{i-1}^+ \pi_{i-1} + T_{i+1}^- \pi_{i+1} - [T_i^+ + T_i^-] \pi_i$$

is solved using the **WKB approach**

Anti-coordination Games & WKB Theory (II)

To solve the QSME $T_{i-1}^+ \pi_{i-1} + T_{i+1}^- \pi_{i+1} - [T_i^+ + T_i^-] \pi_i = 0$ away from the boundaries, one uses the WKB Ansatz ($x = i/N$):

$$\pi(x) = \mathcal{A} e^{-NS(x) - S_1(x)}$$

$S(x)$ is the “action” and $S_1(x)$ is the amplitude, while \mathcal{A} is a constant. With this ansatz and $\mathcal{T}_\pm(x) \equiv T_i^\pm$, one obtains to order $\mathcal{O}(N^{-1})$

$$\begin{aligned} \pi(x) & \left\{ \mathcal{T}_+(x) \left[e^{S'} \left(1 - \frac{1}{2N} S'' + \frac{1}{N} S_1' \right) - 1 \right] \right. \\ & + \mathcal{T}_-(x) \left[e^{-S'} \left(1 - \frac{1}{2N} S'' - \frac{1}{N} S_1' \right) - 1 \right] \\ & \left. + \frac{1}{N} \left[e^{-S'} \mathcal{T}_-'(x) - e^{S'} \mathcal{T}_+'(x) \right] \right\} = 0. \end{aligned}$$

To order $\mathcal{O}(1)$, with the “momentum” $p(x) = dS/dx$:

Hamilton-Jacobi equation, $H[x, S'(x)] = 0$, where the Hamiltonian is $H(x, p) = \mathcal{T}_+(x)(e^p - 1) + \mathcal{T}_-(x)(e^{-p} - 1)$

Anti-coordination Games & WKB Theory (III)

To solve the QSME $T_{i-1}^+ \pi_{i-1} + T_{i+1}^- \pi_{i+1} - [T_i^+ + T_i^-] \pi_i = 0$ away from the boundaries, one uses the WKB Ansatz ($x = i/N$):

$$\pi(x) = \mathcal{A} e^{-NS(x) - S_1(x)}$$

To order $\mathcal{O}(1)$: zero-energy trajectories of Hamiltonian $H[x, S'(x)]$ yields $p_a(x) = -\ln[\mathcal{T}_+(x)/\mathcal{T}_-(x)] \Rightarrow$ “optimal path” to fixation is $S(x) = -\int^x \ln[\mathcal{T}_+(\xi)/\mathcal{T}_-(\xi)] d\xi$

To order $\mathcal{O}(N^{-1})$: $S_1(x)$ by solving a differential equation

Constant \mathcal{A} : by Gaussian normalization of the QSD $\pi(x)$ about x^*

- *To order $\mathcal{O}(1)$:* $S(x) = -\int^x \ln[\mathcal{T}_+(\xi)/\mathcal{T}_-(\xi)] d\xi$
- *To order $\mathcal{O}(N^{-1})$:* $S_1(x) = \frac{1}{2} \ln[\mathcal{T}_+(x)\mathcal{T}_-(x)]$

Near the boundary $x = 0$, expand $\mathcal{T}_\pm(x) \simeq x \mathcal{T}'_\pm(0)$ in the QSME $\Rightarrow \mathcal{T}'_+(0)(i-1)\pi_{i-1} + \mathcal{T}'_-(0)(i+1)\pi_{i+1} - i[\mathcal{T}'_+(0) + \mathcal{T}'_-(0)]\pi_i = 0$,

yielding $\pi_i = \frac{(R_0^i - 1)\pi_1}{(R_0 - 1)^i}$ with $R_0 \equiv \mathcal{T}'_+(0)/\mathcal{T}'_-(0)$.

Similarly with the boundary $x = 1$

Anti-coordination Games & WKB Theory (IV)

WKB solution for the QSD in the bulk (for $N^{-1/2} \ll x \ll 1 - N^{-1/2}$):

$$\pi(x) = \mathcal{I}_+(x^*) \sqrt{\frac{S''(x^*)}{2\pi N \mathcal{I}_+(x) \mathcal{I}_-(x)}} e^{-N[S(x) - S(x^*)]},$$

Near the boundaries, matching the recursive and WKB solutions yields (with $R_1 \equiv \mathcal{I}'_-(1)/\mathcal{I}'_+(1)$):

$$\pi_1 = \sqrt{\frac{NS''(x^*)}{2\pi}} \frac{\mathcal{I}_+(x^*) (R_0 - 1)}{\sqrt{\mathcal{I}'_+(0) \mathcal{I}'_-(0)}} e^{-N[S(0) - S(x^*)]}$$
$$\pi_{N-1} = \sqrt{\frac{NS''(x^*)}{2\pi}} \frac{\mathcal{I}_+(x^*) (R_1 - 1)}{\sqrt{\mathcal{I}'_+(1) \mathcal{I}'_-(1)}} e^{-N[S(1) - S(x^*)]}$$

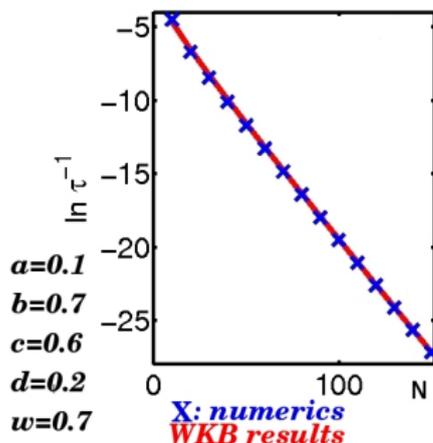
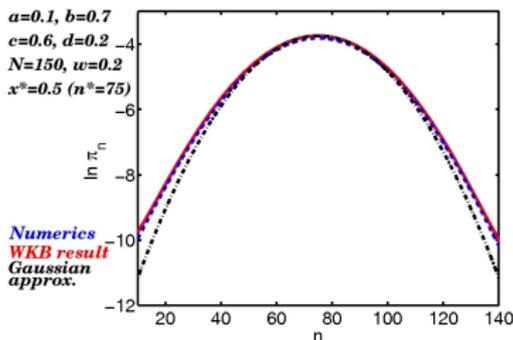
Thus, $\tau = N [\mathcal{I}'_-(0)\pi_1 + |\mathcal{I}'_+(1)|\pi_{N-1}]^{-1}$ and $\phi = N \mathcal{I}'_-(0)\pi_1 \tau$
For the fMP:

$$e^{-NS(x)} = [Ax + B(1-x)]^{Nx - N(\frac{B}{B-A})} [Cx + D(1-x)]^{-Nx - N(\frac{D}{C-D})},$$

with $A = 1 - w + wa$, $B = 1 - w + wb$, $C = 1 - w + wc$, and $D = 1 - w + wd$.

Anti-coordination Games & WKB Theory: Results (I)

- **QSD**: bell-shaped function peaked at x^* . Systematic non-Gaussian effects near the tails, well accounted by the WKB approach
- **MFTs**: exponential dependence on the population size ($Nw \gg 1$), $\tau \propto N^{1/2} e^{N(\Sigma - S(x^*))}$, where $\Sigma \equiv \min(S(0), S(1))$
For “small” selection intensity, the MFTs grow exponentially as $\tau^A \sim N^{1/2} e^{Nw(a-c)^2/[2(c-a+b-d)]}$, $\tau^B \sim N^{1/2} e^{Nw(b-d)^2/[2(c-a+b-d)]}$, and $\tau = \tau^A \tau^B / (\tau^A + \tau^B)$

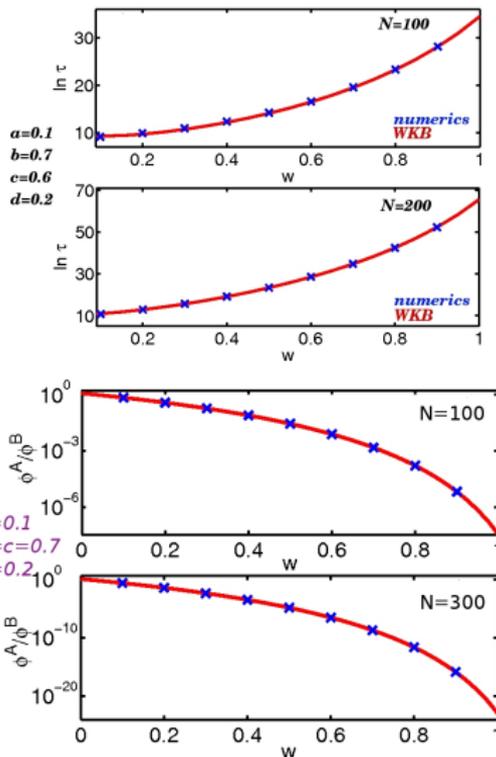


Anti-coordination Games & WKB Theory: Results (II)

- For $Nw \gg 1$, the MFTS increase monotonically with w , faster than exponentially
- *Fixation probability*: When $w = 0$, $\phi^A/\phi^B = x/(1-x)$ depends on initial fraction of mutants. **No longer the case when $w > 0$ (selection):**

$$\frac{\phi^A}{\phi^B} \rightarrow \sqrt{\frac{BD}{AC}} \left(\frac{C-A}{B-D} \right) \times \frac{B^{N(\frac{B}{B-A})} D^{N(\frac{D}{C-D})}}{A^{N(\frac{A}{B-A})} C^{N(\frac{C}{C-D})}}$$

\Rightarrow Exponential dependence: ϕ_A/ϕ_B is exponentially large/small when $N \gg 1$, except for $w \ll 1$



In CGs, $i = 0$ and $i = N$ are attractors and x^* is unstable.
Starting with i A individuals, what is the probability ϕ_i^A that species A fixates the system?

ϕ_i^A is a cumulative distribution function obeying

$$T_i^+ \phi_{i+1}^A + T_i^- \phi_{i-1}^A - [T_i^+ + T_i^-] \phi_i^A = 0, \quad \text{with} \quad \phi_0^A = 0, \phi_N^A = 1$$

Convenient to work with $\mathcal{P}_i = \phi_{i+1}^A - \phi_i^A$ such that $\phi_i^A = \sum_{m=0}^{i-1} \mathcal{P}_m$.
When $N \gg 1$, $\mathcal{P}_i = \mathcal{P}(x)$ and the latter obeys

$$\mathcal{I}_+(x) \mathcal{P}(x) - \mathcal{I}_-(x) \mathcal{P}(x - N^{-1}) = 0.$$

Eq. solved by the WKB ansatz

$$\mathcal{P}(x) = \mathcal{A}_{CG} e^{-N\mathcal{S}(x) - \mathcal{I}_1(x)}$$

As for ACGs, this leads to $\mathcal{S}(x) = -S(x) = \int^x \ln[\mathcal{I}_+(\xi)/\mathcal{I}_-(\xi)] d\xi$
and $\mathcal{I}_1(x) = -\frac{1}{2} \ln[\mathcal{I}_-(x)/\mathcal{I}_+(x)]$

One therefore obtains:

$$\mathcal{P}(x) = \sqrt{\frac{|S''(x^*)|}{2\pi N} \frac{\mathcal{I}_-(x)}{\mathcal{I}_+(x)}} e^{N[S(x) - S(x^*)]}$$

To leading order when $N^{-1} \ll w \ll 1$:

$$\phi^A(x) \simeq \sqrt{\frac{N|S''(x^*)|}{2\pi}} \int_0^x dy e^{N[S(y) - S(x^*)]}$$

Criterion of evolutionary stability (of “wild species” B): $\phi^A(x) < x$, for $x \ll 1 \Rightarrow$ *relevant to consider the limit $x \ll x^*$ with finite w*

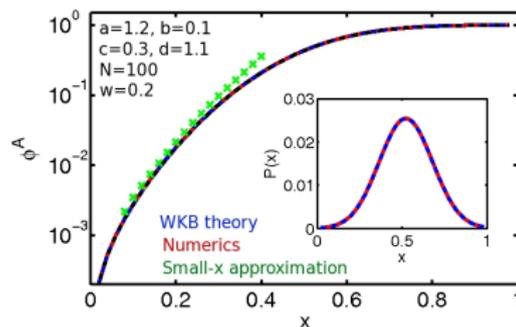
Approximation for $N^{-1} \ll x \ll 1$ (where $S'(x) > 0$) and $Nw \gg 1$:

$$\phi^A(x) \simeq \frac{\mathcal{P}(x)}{e^{S'(x)} - 1}$$

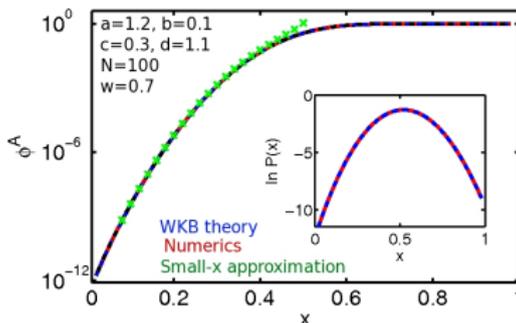
As $\phi^A(x)$ is exponentially small, $\phi^A(x) < x$ and the selection opposes replacement of B 's by A 's \Rightarrow the state with all B 's is always evolutionary stable when w is finite

Coordination Games & WKB Theory: Results (I)

- *Fixation probability:* $\phi^A(x) \rightarrow 1$ when $x \rightarrow 1$, with $\phi^A(x^*) = 1/2$, and is exponentially small $\phi^A \rightarrow 0$ when $x \rightarrow 0$. “Jump” from finite to exponentially small value of ϕ^A becomes steeper when w increases



- *Behaviour for $x \ll 1$:* When w is finite, $N \gg 1$ and $x \ll 1$, the exponentially small value of $\phi^A(x)$ is approximated by
$$\phi^A(x) \simeq \frac{\mathcal{P}(x)}{e^{S'(x)-1}}$$



Coordination Games & WKB Theory: Results (II)

Comparison with Fokker-Planck:

Fixation probability often approximated using the Fokker-Planck Equation (FPE).

This diffusion approx. yields

$$\phi_{\text{FPE}}^A(x) = \frac{\Psi(x)}{\Psi(1)} \text{ with}$$

$$\Psi(x) = \int_0^x e^{-\int_0^y \Theta_{\text{FPE}}(z) dz} dy \text{ and}$$

$$\Theta_{\text{FPE}}(x) = 2N \left(\frac{\mathcal{F}_+(x) - \mathcal{F}_-(x)}{\mathcal{F}_+(x) + \mathcal{F}_-(x)} \right)$$

Often used within linear noise approx., where $\phi_{\text{FPE}}^A(x) = \frac{\Psi(x)}{\Psi(1)}$ with

$$\Theta_{\text{FPE}}(x) =$$

$$2N(x - x^*) \left(\frac{\mathcal{F}'_+(x^*) - \mathcal{F}'_-(x^*)}{\mathcal{F}_+(x^*) + \mathcal{F}_-(x^*)} \right) \text{ instead}$$

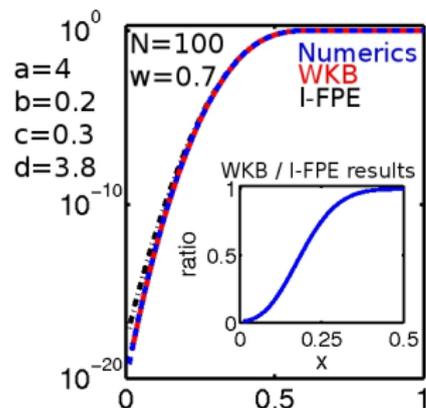
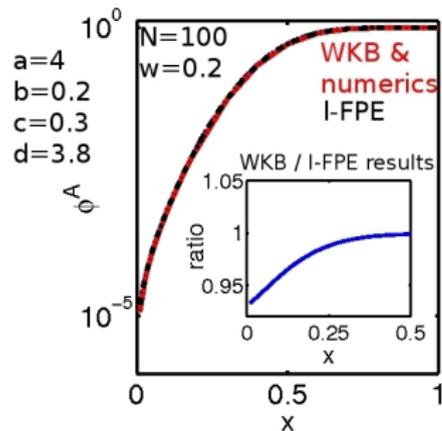
of $\Theta_{\text{FPE}}(x)$

To leading order, WKB result can

be rewritten as $\phi^A(x) \simeq \frac{\Psi(x)}{\Psi(1)}$, with

$\Theta(x) = N \ln[\mathcal{F}_+(x)/\mathcal{F}_-(x)]$ instead

of $\Theta_{\text{FPE}}(x)$

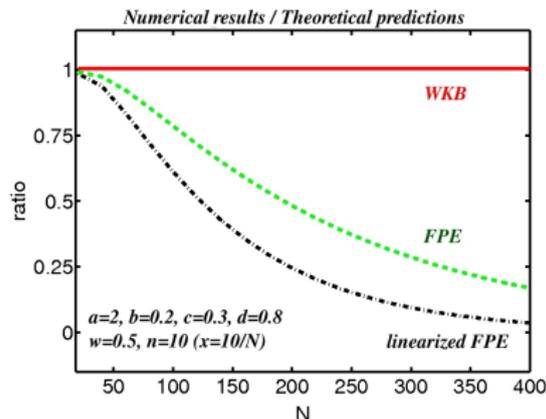
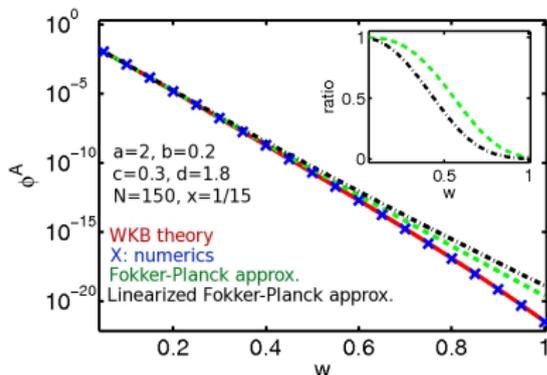


Coordination Games & WKB Theory: Results (III)

- Excellent agreement between numerics and WKB results for any x and $w > 0$
- FPE in good agreement with WKB and numerics when w is small (and/or $x \simeq x^*$).
- However, exponentially large deviations when w and N are raised and x deviates from x^*

As $\Theta(x) - \Theta_{\text{FPE}}(x) \sim N(w\Delta x)^3$
 and $\Theta(x) - \Theta_{\text{LFPE}}(x) \sim N(w\Delta x)^2$
 ($\Delta x = x - x^*$) \Rightarrow

Exponentially large errors in $\phi_{\text{FPE}}^A(x)$ and $\phi_{\text{LFPE}}^A(x)$ when $w \gtrsim N^{-1/3}$ and $w \gtrsim N^{-1/2}$, resp.



Presentation of a WKB-based approach allowing to compute large-fluctuation-induced phenomena in evolutionary processes

- Account naturally for large fluctuations and non-Gaussian behaviour
- Application to a class of evolutionary games modelling: *combined effect of stochasticity and non-linearity (selection)?*
- Metastability in Anti-Coordination Games: calculation of the QSD, ϕ and MFTs \Rightarrow when $w > 0$ and $N \gg 1$, non-Gaussian QSD and MFTs grow exponentially with N
- ϕ^A in Coordination Games: asymptotically exact results for $\phi^A \Rightarrow$ exponentially small when $w > 0$ and $N \gg 1$
- Comparison with Fokker-Planck: FPE is only accurate around x^* and for vanishingly small selection strength w
- Generalization to other rules/interactions
- Method can be adapted to study non-exactly solvable problems (e.g. 3×3 games)